

Chemical Reaction and Thermal Radiation Effects on Magnetohydrodynamic Nanofluid Flow Past an Exponentially Stretching Sheet

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Abstract

This study investigates chemical reaction and thermal radiation effects on hydromagnetic nanofluid flow over an exponentially stretching sheet. The governing partial differential equations were transformed to ordinary differential equations by using similarity transformation and the resulting equations were solved using asymptotic series method. Graphical results showing the influence of the governing parameters on the velocity, temperature and concentration are displayed. Our results indicated that an increase in stretching sheet, thermal Grashof parameters leads to the increase in the rate of fluid flow while it decreases when magnetic field factor is increased. Also, increasing the thermophoresis number brings about increase in temperature and concentration while the reverse is the case as Prandtl number, Schmidt factor and chemical reaction rate increases. Increase in radiation leads to increase in the temperature.

Keywords: Thermophoresis, Stretching sheet, MHD (Magnetohydrodynamic), thermal Radiation, Chemical reaction, Nanofluid.

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1. Introduction

The study of magnetohydrodynamic nanofluid flow over an exponentially stretching sheet exhibits several healthy applications in science and technology, industrial appliances, electronic gadgets, automobile industries, etc. However, the areas of application and usefulness of this study is vast and include nuclear reactor fluidization: In order to have effective cooling process in the reactors of nuclear fission, electromagnetic forces are used to supply liquified sodium around the reactors, cooling of an infinite hot metal sheet in a cooling tub, rubber sheets and plastic production, drawing of cables and fiber sheets, etc. Thermal as well as mass transmission in Carreau fluid broadens the importance and useful applications in sciences such as physiology, symptoms of neurology, treatment of diagnostic diseases etc. However, it is important to mention that although nanofluid (Nano suspension) is made up of two parts which are the base fluid and solid particles, “surfactant” is added in some cases in order to elongate the stability of the mixture. (Sheremet and Pop, 2015), investigated steady-state free convection in a medium containing pore having nanofluid, couple energy transference and Buongiorno’s model. They noted that the local Nusselt number rises with a rise in Rayleigh number and buoyancy ratio index.

According to (Chen, 1998) an investigation was carried on the motion of fluid and thermal transmission over a surface that is elongated was done. A modification on this study was carried out as reported by (Ishak *et al.*, 2008). It was noted that such physical parameters of interest decline when magnetic force/strength rises.

The analysis of magnetohydrodynamic Casson fluid over a stretching surface past spongy material indicated that the rate of fluid movement retards because of the drag force produced by transverse magnetic source (Nadeem *et al.*, 2013).

The study of flow over boundary layer past a strong sheet in motion in the presence of non-varying velocity was championed by (Sakiadis, 1961). His work was revisited through the introduction of a extending material in a linear motion from a non-movable position as given by (Crane, 1970).

Techniques to control flow means a little deviation representing perfectly a great application in field of Engineering such as haul decrease, hoist rise, blending upliftment (Wikipedia, the free encyclopedia, 2017).

However, (Sheremet and Pop, 2014), carried out study on the two-dimensional non-turbulence free convection of water-based fluids in a double porous hollow such that lateral parapets possess sinusoidal parapet degree of coldness or hotness by using Buongiorno’s model. They resolved from their findings that with a slight Lewis number is slight, and increased value of thermophoresis, the heterogeneous smaller particles dispersion in the fluid will be more physical.

Analysis of the impact of different nanoparticles on transmission of heat as a result of impartation of temperature gradient over a thermally medium. It was reported that transmission of heat energy is well pronounced (Oztop and Abu-Nada, 2008). (Santoshi and Govardhan, 2020), considered the influence of magnetic strength, heat source/sink in nanofluid past a permeable elongating surface with external heat

energy and non-complete slide. (Rohana, 2014) examined the magnetohydrodynamic Marangonic convective nanofluid in the presence of injection and blowing. Verification of magnetohydrodynamics boundary layer passage of thermal cum mass transmission past upright material in motion, chemically reacting in presence of suction by (Ibrahim and Makinde, 2010) was noted. Indications from their findings depicts that rate of fluid motion suppresses as temperature and concentration elevates.

Studies abound on the influence of thermal radiation on fluid flows and the results obtained have been significant, (Raptis *et al.*, 2004), (Mabood *et al.*, 2017), (Devi and Reddy, 2014), (Bunonyo *et al.*, 2018) and (Omamoke *et al.*, 2020) just to mention but a few. However, influence of radiation, viscosity on MHD non-steady natural flow past a partially sieve-like material was analyzed by (Joaquin, 2007). The influence of radiative temperature on nanofluid movements instigated via a radically overextended exterior through flexible viscidness was reported by (Makinde *et al.*, 2016) and they found that resistance issue appreciates by means of convective constraint. (Raju *et al.*, 2016), categorized the impacts of hotness and mass conveyance on magnetohydrodynamic movement past an enlarging surface. From their study, exponential factor has tendency to elevate the degree of fluid movement. Meanwhile, (Irfan *et al.*, 2019) considered the magnetohydrodynamic free stream and heat transfer of nanofluid flow over an exponentially radiating stretching sheet with variable fluid properties. They opined that thermal boundary coating stiffness rises as the thermal conductivity, viscosity of fluid, Brownian motion and Thermophoresis numbers intensify.

However, as a result of the numerous applications of nanofluid flow over a stretching surface, many researchers and scientists have been attracted to its study. It's also imperative to remark that magnetohydrodynamic fluids possess the characteristics of electromagnetic conductivity which is very important as long as flow and heat transfer is concerned. Nevertheless, the application of this concept in areas of modern technological advancement has really been shown to be very useful to humans. For example, in the satellite and solar power industries, medical fields, etc. with regards to the medical field, the nanoparticles made of gold are used in the manufacturing and distribution of medicine for the treatment of tumors of cancer in the body (Bouslimi *et al.*, 2021).

(Dandpat and Chakraborty, 2010), analyzed the outcome of changeable fluid factors on tinny liquid motion picture flow over an instable intense widening sheet. They noted that the consequence of inconstant fluid thickness on the profile of the velocity enhances as viscosity decreases as a result of temperature decrease on the stretching plate. (Hsiao, 2001), investigated the magneto-hydrodynamic inaction point passage of visco-elastic fluid on thermal making stretching sheet with the effect of viscid degeneracy. They opined that as the magnetic parameter upsurges, the proportion of assignment of heat energy falls. However, analysis of the flow over a shaky elongating surface with chemical reaction and unequal heat origin by applying Runge-Kutta-Fehlberg technique and shooting approaches was carried out by (Seini, 2013). It was indicated that the amount of transmission of heat and mass

together with the coefficient of the skin friction enhances with an enhancement of the unsteadiness factor but retards with an enhancement of the space-dependent and temperature-dependent factors.

Similarly, an investigation of magneto-hydrodynamic limit layer flow because of exponential stretching surface through radiation and chemical reaction by applying Runge-Kutta-Fehlberg technique and shooting approaches was executed by (Seini and Makinde, 2013). As a result of the magnetic parameter, the presence of a velocity term which appears nonlinear was found in the equation of the temperature of the model. Their results showed that the enhancement of the parameters of radiation as well as transverse magnetic field brings about to decrease in the transfer heat energy proportion on the surface. In the same vein, (Sahoo and Biswal, 2015), presented magneto-hydrodynamic viscoelastic boundary layer flow over a stretching surface in the presence of heat transfer. They indicated that increased values of Prandtl factor causes a gradual decay of temperature because of low thermal conductivity.

From the studies examined above, none of them investigated the influence of chemical reaction on the flow. Hence, we carry out our analysis on an exponentially stretching sheet in a hydromagnetic nanofluid, putting into consideration stretching material, convective terms, rate of chemical reaction, thermal conductivity and radiation.

2. Mathematical Formulation

We consider a steady dual-dimensional, electrically conducting, viscous, incompressible flow of nanofluid over a stretching sheet. The stretching sheet is along the x -axis while the y -axis is normal to it. The velocity components are u, v along the x -, y -, axis respectively. The flow field is exposed to a magnetic field of strength B_0 in the path transverse to the surface of the stretching surface. We assume that the magnetic Reynolds number is small, hence the induced magnetic field is negligible and consequently considered only the applied magnetic field. Meanwhile, the nanoparticles immersed in the base fluid are influenced as a result of Brownian motion. The temperature, T and nanoparticles volume fraction C at the boundary ($y = 0$) are T_w and C_w respectively while far from the plate the temperature and nanoparticles fraction are T_∞ and C_∞ respectively. Therefore, the boundary layer equations governing the present study are presented as follows:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0 \quad (1)$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = u_\infty \frac{du_\infty}{dx} + \frac{1}{\rho} \frac{\partial}{\partial y} \left(\mu \frac{\partial u}{\partial y} \right) - \frac{\sigma B^2}{\rho} (u - u_\infty) + \frac{g\beta_T l^3 (T - T_\infty)}{v^2} \quad (2)$$

$$u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} = \frac{1}{\rho c_p} \frac{\partial}{\partial y} \left(\frac{k \partial T}{\partial y} \right) + \tau \frac{D_B(C-C_\infty)}{D_T T_\infty} - \frac{1}{\rho c_p} \frac{\partial q_r}{\partial y} \quad (3)$$

$$u \frac{\partial C}{\partial x} + v \frac{\partial C}{\partial y} = D_B \frac{\partial^2 C}{\partial y^2} + \frac{D_T}{T_\infty} (T - T_\infty) - k_r (C - C_\infty) \quad (4)$$

Subjected to the following boundary conditions:

$$\left. \begin{aligned} u = U_w(x) = a e^{\frac{x}{2L}}, v = -v_w, T = T_w, C = C_w \text{ at } y = 0 \\ u \rightarrow U_\infty(x) = b e^{\frac{x}{2L}}, T \rightarrow T_\infty, C \rightarrow C_\infty \text{ as } y \rightarrow \infty \end{aligned} \right\} \quad (5)$$

The similarity transformation variables are defined as follows:

$$\eta = \sqrt{\frac{a}{2vL}} e^{\frac{x}{2L}} y, \varphi = \sqrt{2vLa} e^{\frac{x}{2L}} f(\eta), \theta = \frac{T-T_\infty}{T_w-T_\infty}, \quad \phi = \frac{C-C_\infty}{C_w-C_\infty}, \quad (6)$$

$$u = \frac{\partial \psi}{\partial y}, \quad v = -\frac{\partial \psi}{\partial x}$$

where, ρ_p is the nanoparticle density, ρ_f is the base fluid density, β is volumetric thermal expansion coefficient of the base fluid, σ is electrical conductivity parameter, g is the acceleration due to gravity, k is thermal conductivity, μ is viscosity of the fluid, D_T is thermophoresis diffusion parameter, k_r is the chemical reaction factor.

3. Method of solution

In order to solve the problem, we shall use the variables stated in equation (6) to transform equations (1) – (4) to have:

$$f''''(\eta) + 2(\lambda^2 - (f')^2(\eta)) + f(\eta)f''(\eta) - M(f'(\eta) - \lambda) + Gr\theta(\eta) = 0 \quad (7)$$

$$\left(1 + \frac{4R}{3}\right)\theta''(\eta) + Pr[(f(\eta)\theta'(\eta) + f'(\eta)\theta(\eta))] + \frac{N_b}{N_t}W(\eta) = 0 \quad (8)$$

$$W''(\eta) + \frac{N_t}{N_b}\theta(\eta) + Sc(f(\eta)W'(\eta) - f'(\eta)W(\eta) - \gamma ScW(\eta)) = 0 \quad (9)$$

Similarly, the transformed boundary conditions for the second model are:

$$\left. \begin{aligned} f(0) = h_0, f'(0) = 1, \theta(0) = 1, \phi(0) = 1 \text{ at } \eta = 0 \\ f'(\infty) = \lambda, \theta(\infty) \rightarrow 0, \phi(\infty) \rightarrow 0 \text{ as } \eta \rightarrow \infty \end{aligned} \right\} \quad (10)$$

where, $M = \frac{2\sigma B^2 L}{\rho U_w}$, magnetic parameter, $\lambda = \frac{b}{a}$, stretching sheet factor, $Gr = \frac{g\beta_T l^3 (T_w - T_\infty)}{(U_w \nu)^2}$, is the thermal Grashof number, $Pr = \frac{\mu c_p}{k}$, is the Prandtl number, $\tau = \frac{(\rho c_p)_p}{(\rho c_p)_f}$, is ratio of the specific heat capacities of the nanoparticle to that of the base fluid, $R = \frac{4T_\infty^3 \sigma^*}{3k^* k}$, radiation parameter, $N_b = \frac{2\tau D_B (C_w - C_\infty)}{k U_w}$, is Brownian motion factor, $N_t = \frac{2\tau D_T (T_w - T_\infty)}{\nu T_\infty l}$, is thermophoresis factor, $Sc = \frac{\nu}{D_B}$ is Schmidt number, $\gamma = \frac{2Lkr}{U_w}$ = modified chemical reaction rate, and the primes represent differentiation with respect to η , $h_0 = \frac{\nu_w}{\sqrt{\frac{av}{2L}}}$ is the suction parameter.

In order to solve equations (7)-(9) with boundary equation (10), we follow (Bestman, 1990), thus,

$$\eta = \Omega h_0, h(\eta) = h_0 H(\eta), \theta(\eta) = \Theta(\eta), w(\eta) = \Psi(\eta), \varepsilon = \frac{1}{h_0^2} \quad (11)$$

and the resulting set of equations are

$$H'''(\eta) + 2\lambda^2 \varepsilon^2 - 2(H')^2(\eta) + H(\eta)H''(\eta) - MH'(\eta)\varepsilon + M\lambda\varepsilon^2 + \varepsilon^2 Gr\theta(\eta) = 0 \quad (12)$$

$$\theta''(\eta) + \theta''(\eta)\frac{4}{3}R + PrH(\eta)\theta'(\eta) + PrH'(\eta)\theta(\eta) + \varepsilon\frac{N_t}{N_b}\Psi(\eta) = 0 \quad (13)$$

$$\Psi''(\eta) + \varepsilon\frac{N_t}{N_b}\Theta(\eta) + H(\eta)\Psi'(\eta)Sc - H'(\eta)\Psi(\eta)Sc - \varepsilon\gamma Sc\Psi(\eta) = 0 \quad (14)$$

Subject to:

$$H(0) = 1, H'(0) = \varepsilon, H'(\infty) = \varepsilon\lambda, \theta(0) = 1, \theta(\infty) = 0, \Psi(0) = 1, \Psi(\infty) = 0 \quad (15)$$

With $\varepsilon \ll 1$, we shall solve the equation by adopting the asymptotic series technique, Bestman (1990):

$$H(\eta) = 1 + \varepsilon H_1(\eta) + \varepsilon^2 H_2(\eta) + \dots \quad (16)$$

$$\theta(\eta) = \theta_0(\eta) + \varepsilon \theta_1(\eta) + \dots \quad (17)$$

$$\Psi(\eta) = \Psi_0(\eta) + \varepsilon \Psi_1(\eta) + \dots \quad (18)$$

Using equations (12) – (14) on equations (16) – (18) and after differentiations and simplification we obtain the series of approximations as:

$$\left(1 + \frac{4}{3}R\right)\theta_0''(\eta) + Pr\theta(\eta) = 0 \quad (19)$$

$$\Psi_0''(\eta) + Sc\Psi_1'(\eta) = 0 \quad (20)$$

$$H_1'''(\eta) + H_1''(\eta) = 0 \quad (21)$$

$$\left(1 + \frac{4}{3}R\right)\theta_1''(\eta) + Pr\theta_1'(\eta) + PrH_1(\eta)\theta_0'(\eta) + PrH_1'(\eta)\theta_0(\eta) + \frac{Nt}{Nb}\Psi_0(\eta) = 0 \quad (22)$$

$$\Psi_1''(\eta) + Sc\Psi_1'(\eta) + H_1(\eta)\Psi_0'(\eta)Sc + H_1'(\eta)\Psi_0(\eta)Sc + \frac{Nt}{Nb}\theta_0(\eta) - \gamma Sc\Psi_0(\eta) = 0 \quad (23)$$

$$H_2''(\eta) + 2\lambda^2 - 2(H_1')^2 + (\eta) + H_2''(\eta) + H_1(\eta)H_1''(\eta) - MH_1'(\eta) + M\lambda + Gr\theta(\eta) = 0 \quad (24)$$

Subject to:

$$\theta_0(0) = 1, \theta_0(\infty) = 0, \Psi_0(0) = 1, \Psi_0(\infty) = 0, H_1(0) = 0, H_1'(0) = 1, H_1'(\infty) = \lambda \quad (25)$$

$$\theta_1(0) = 0, \theta_1(\infty) = 0, \Psi_1(0) = 0, \Psi_1(\infty) = 0, H_2(0) = 0, H_2'(0) = 0, H_2'(\infty) = 0$$

Thus, solving equations (19) – (24) subject to equation (25), we have the following solutions:

$$\begin{aligned} h'(\eta) = & (\lambda\eta - \lambda + 1 + \lambda e^{-\eta} - e^{-\eta}) - \left(4(\lambda)^2\eta e^{-\eta} - 4\lambda\eta e^{-\eta} - \lambda e^{-2\eta} + \right. \\ & \left. \frac{(\lambda)^2}{2} e^{-2\eta} + \frac{1}{2} e^{-2\eta} + M\lambda\eta e^{-\eta} - M\eta e^{-\eta} + \frac{C_1\lambda}{2}(\eta)^2 e^{-\eta} + C_1\lambda\eta e^{-\eta} + C_2\lambda\eta e^{-\eta} - \right. \\ & \left. \frac{C_3\lambda}{4} e^{-2\eta} - \frac{C_1}{2}(\eta)^2 e^{-\eta} - C_1\eta e^{-\eta} - C_2\eta e^{-\eta} + \frac{C_3}{4} e^{-2\eta} - \frac{Gr}{q(q-1)} e^{-q\eta} - \frac{(\lambda)^2}{2} + 3\lambda - \right. \\ & \left. M\lambda + M - C_1\lambda - C_2\lambda - \frac{C_3\lambda}{2} + C_1 + C_2 + \frac{C_3}{4} - \frac{Gr}{q-1} + \frac{1}{2} + \frac{Gr}{q(q-1)} - \left(3(\lambda)^2 - 2\lambda - \right. \right. \\ & \left. \left. 1 + M\lambda - M + C_1\lambda + C_2\lambda + \frac{C_3\lambda}{2} - C_1 - C_2 - \frac{C_3}{2} + \frac{Gr}{q-1}\right) e^{-\eta} \right) \end{aligned}$$

Where,

$$C_1 = \lambda$$

$$\begin{aligned}
C_2 &= 1 \\
C_3 &= \lambda - 1 \\
\theta(\eta) &= e^{-q\eta} + \left(-\frac{q\lambda}{2}(\eta)^2 e^{-q\eta} + q\lambda\eta e^{-q\eta} - q\eta e^{-q\eta} + \frac{(q)^2\lambda}{1+q} e^{-(1+q)\eta} - \right. \\
&\quad \left. \frac{(q)^2}{1+q} e^{-(1+q)\eta} + \frac{q\lambda}{1+q} e^{-(1+q)\eta} - \frac{q}{1+q} e^{-(1+q)\eta} - \frac{3Nt}{NbSc(3+4R)(-q+Sc)} e^{-Sc\eta} - \right. \\
&\quad \left. \frac{(q)^2\lambda}{1+q} e^{-q\eta} + \frac{(q)^2}{1+q} e^{-q\eta} - \frac{q\lambda}{1+q} e^{-q\eta} + \frac{q}{1+q} e^{-q\eta} + \frac{3Nt}{NbSc(3+4R)(Sc-q)} e^{-q\eta} \right) \quad (26)
\end{aligned}$$

$$\begin{aligned}
\Psi(\eta) &= e^{-Sc\eta} + \frac{\lambda}{2}(\eta)^2 e^{-Sc\eta} - 2\lambda\eta e^{-Sc\eta} + \eta e^{-Sc\eta} - \frac{Sc\lambda}{1+Sc} e^{-(1+Sc)\eta} + \\
&\quad \frac{Sc}{1+Sc} e^{-(1+Sc)\eta} - \frac{Sc\lambda}{1+Sc} e^{-(1+Sc)\eta} + \frac{Sc}{1+Sc} e^{-(1+Sc)\eta} - \frac{Nt}{Nbq(q-Sc)} e^{-q\eta} - \gamma\eta e^{-Sc\eta} + \\
&\quad \frac{Sc\lambda}{1+Sc} e^{-Sc\eta} - \frac{Sc}{1+Sc} e^{-Sc\eta} + \frac{Sc\lambda}{1+Sc} e^{-Sc\eta} - \frac{Sc}{1+Sc} + \frac{Nt}{Nbq(q-Sc)} e^{-Sc\eta} \quad (27)
\end{aligned}$$

4. Results and Discussion

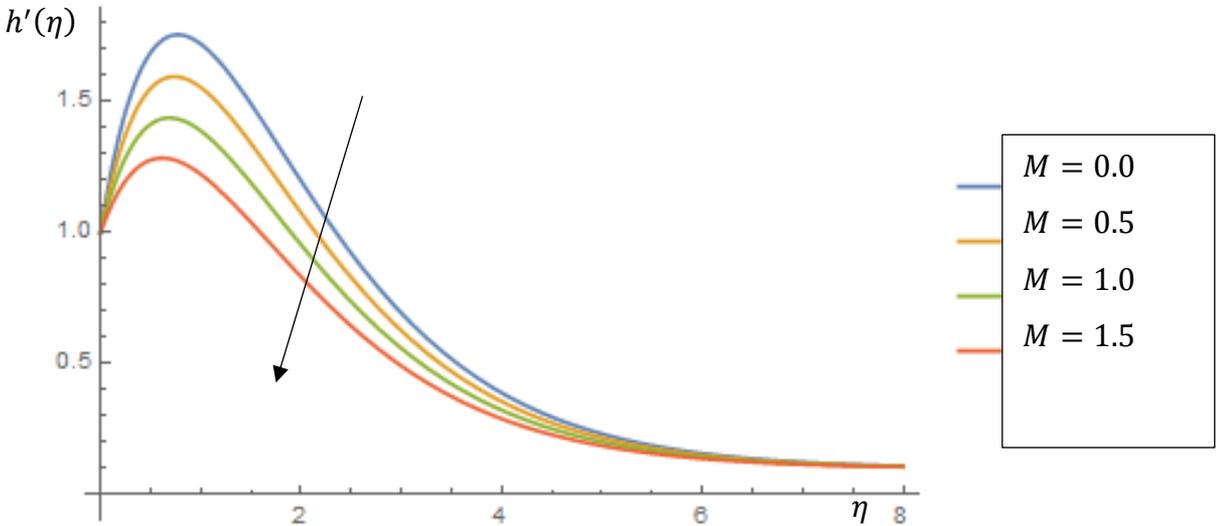


Figure 1: Impact of Magnetic factor M , on Velocity with $\lambda = R = 0.1$, $Gr = 5.0$, $Pr = 1.1$

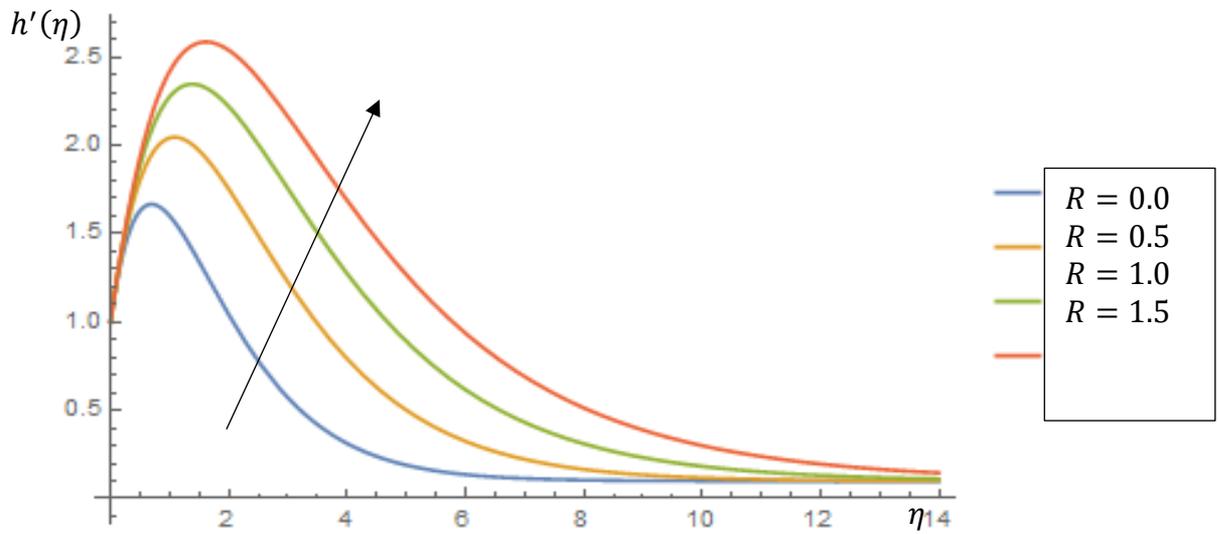


Figure 2: Impact of Radiation R , on Velocity with $\lambda = M = 0.1$, $Gr = 5.0$, $Pr = 1.1$

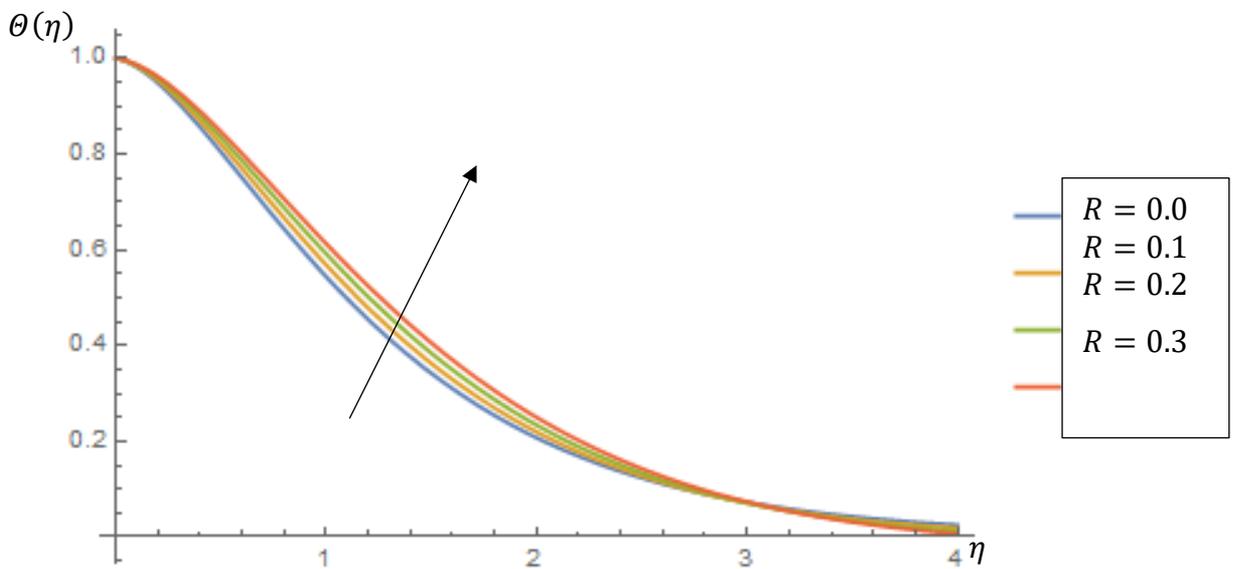


Figure 3: Impact of Radiation R , on Temperature with $Sc = 1.0$, $Nt = Nb = \lambda = 0.1$, $Pr = 1.1$

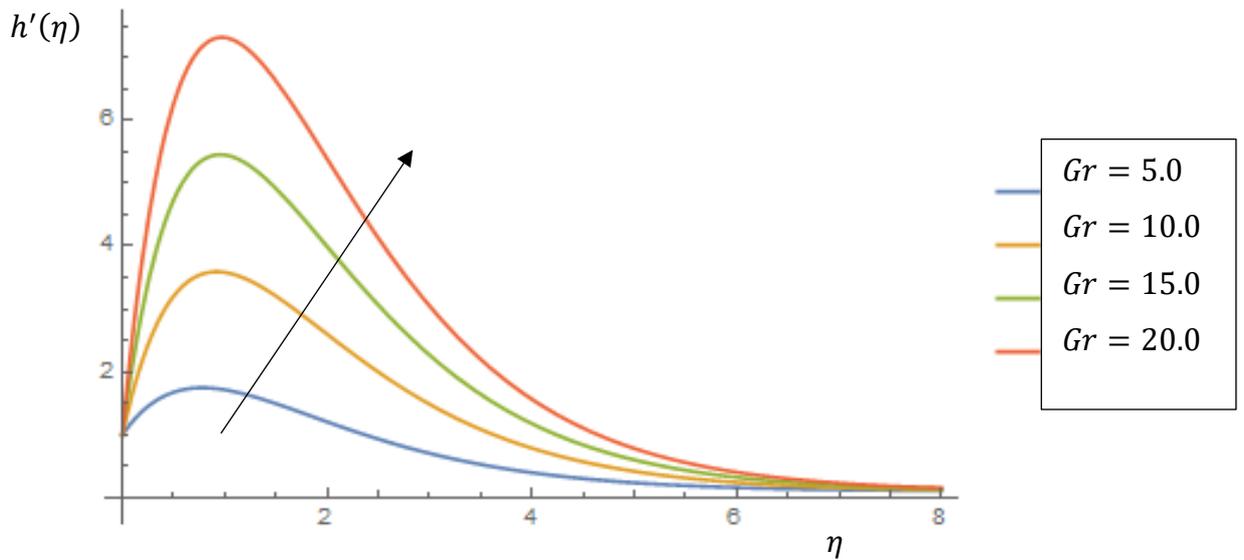


Figure 4: Impact of Modified Grashof number Gr , on Velocity with $\lambda = R = 0.1$, $M = 0.1$, $Pr = 1.1$

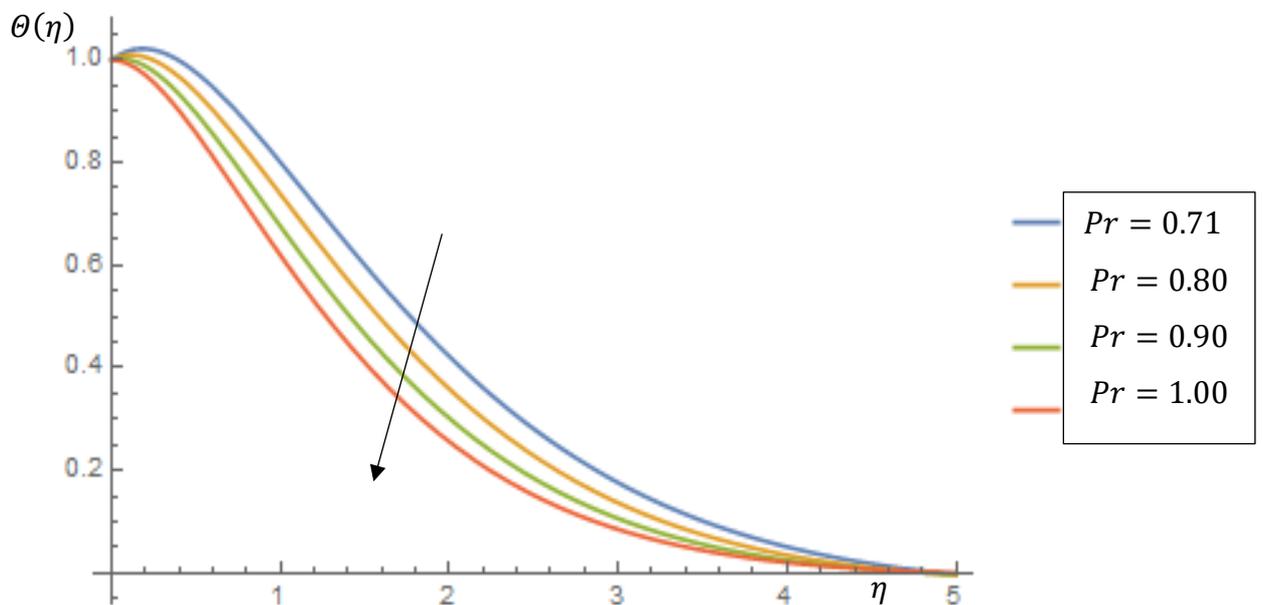


Figure 5: Impact of Prandtl Number Pr , on Temperature with $Sc = 1.0$, $Nt = R = \lambda = Nb = 0.1$

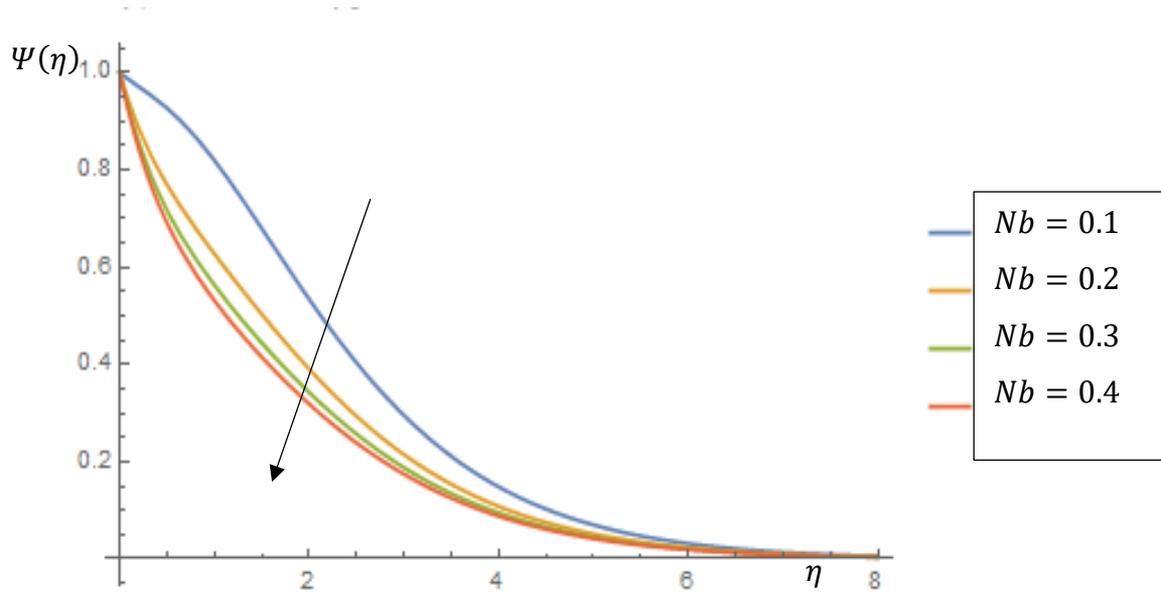


Figure 6: Impact of Brownian motion parameter Nb , on Concentration with $Sc = 1.0$, $Nt = \lambda = \gamma = R = 0.1$, $Pr = 1.1$

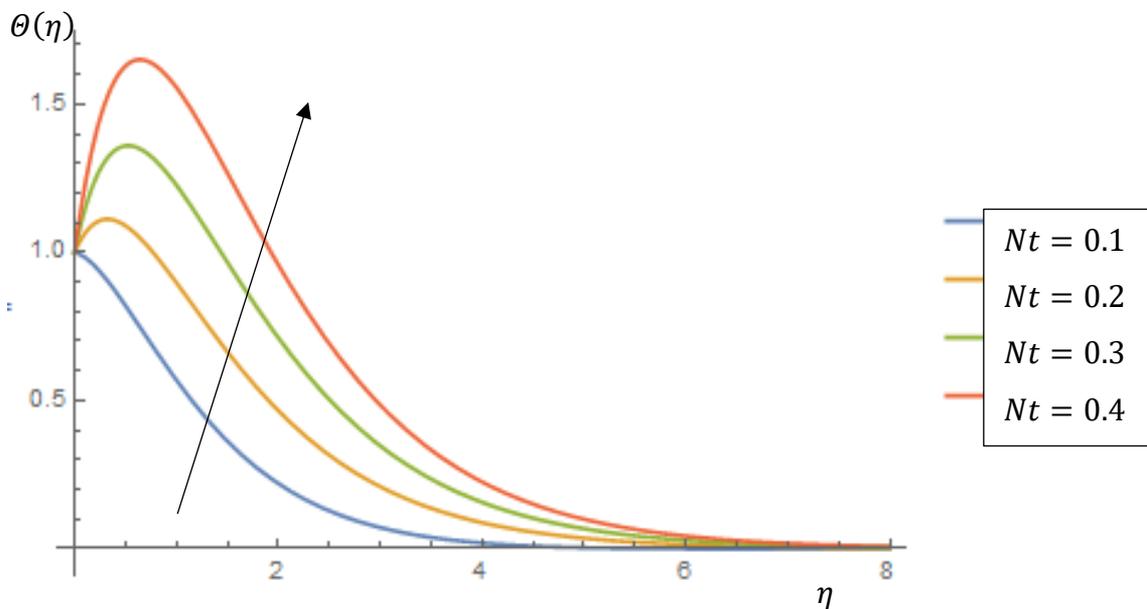


Figure 7: Impact of thermphoresis parameter Nt , on Temperature with $Sc = 1.0$, $Nb = R = \lambda = 0.1$, $Pr = 1.1$

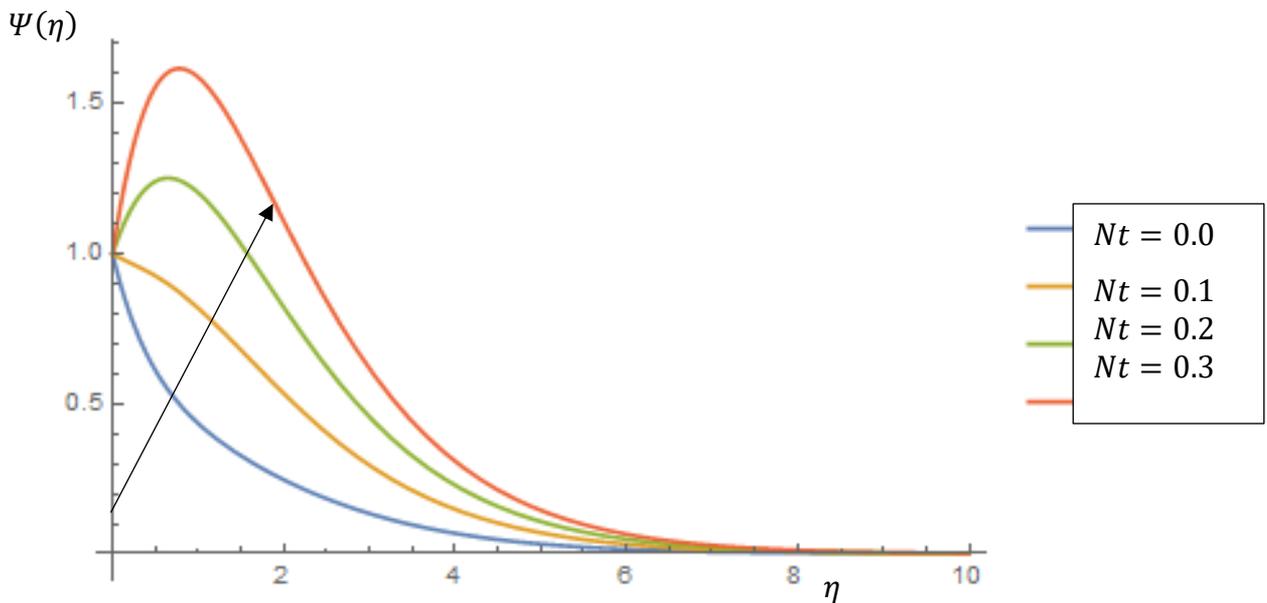


Figure 8: Impact of Thermophoresis parameter Nt , on Concentration with $Sc = 1.0$, $Nb = \lambda = \gamma = R = 0.1$, $Pr = 1.1$

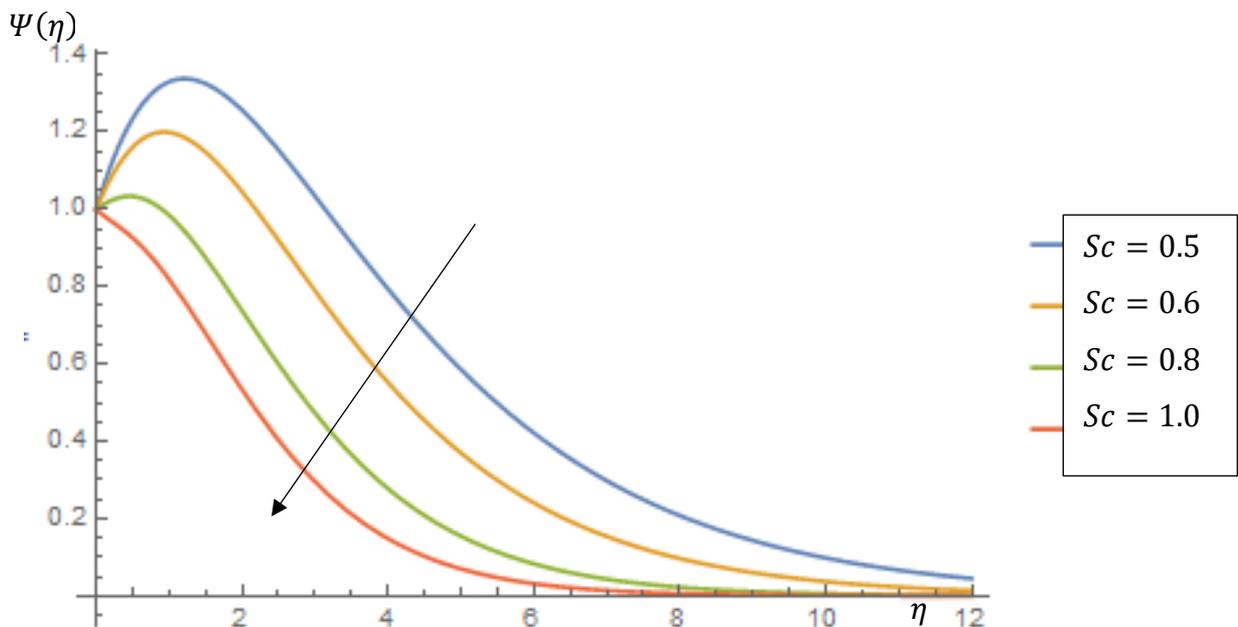


Figure 9: Impact of Schmidt number Sc , on Concentration with $Nb = Nt = \lambda = \gamma = R = 0.1$, $Pr = 1.1$

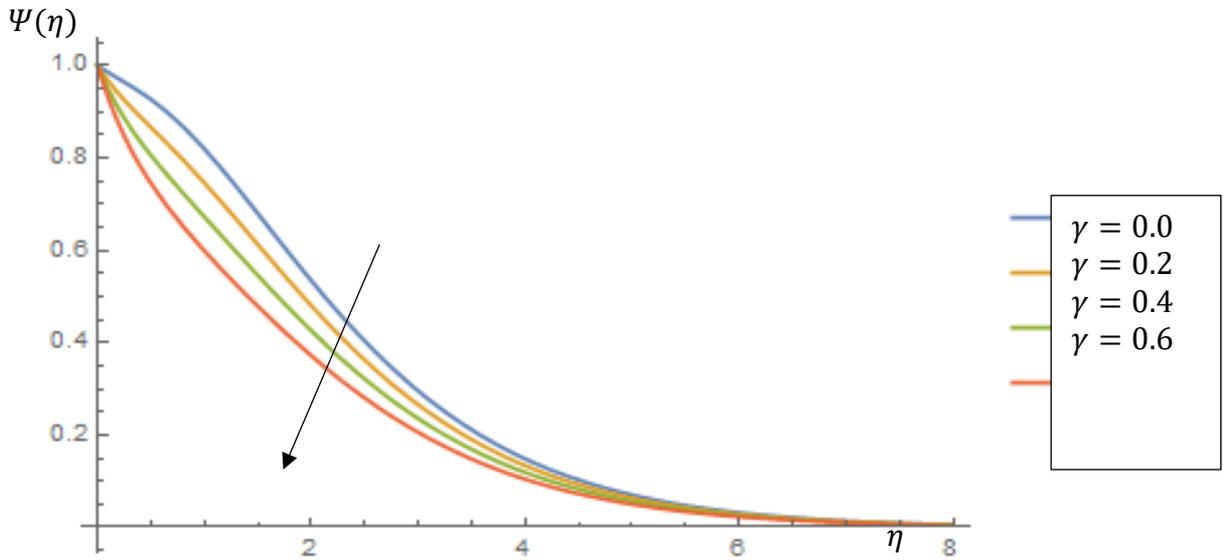


Figure 10: Impact of Rate of Chemical reaction parameter γ , on Concentration with $Nb = Nt = \lambda = R = 0.1, Pr = 1.1, Sc = 1.1$

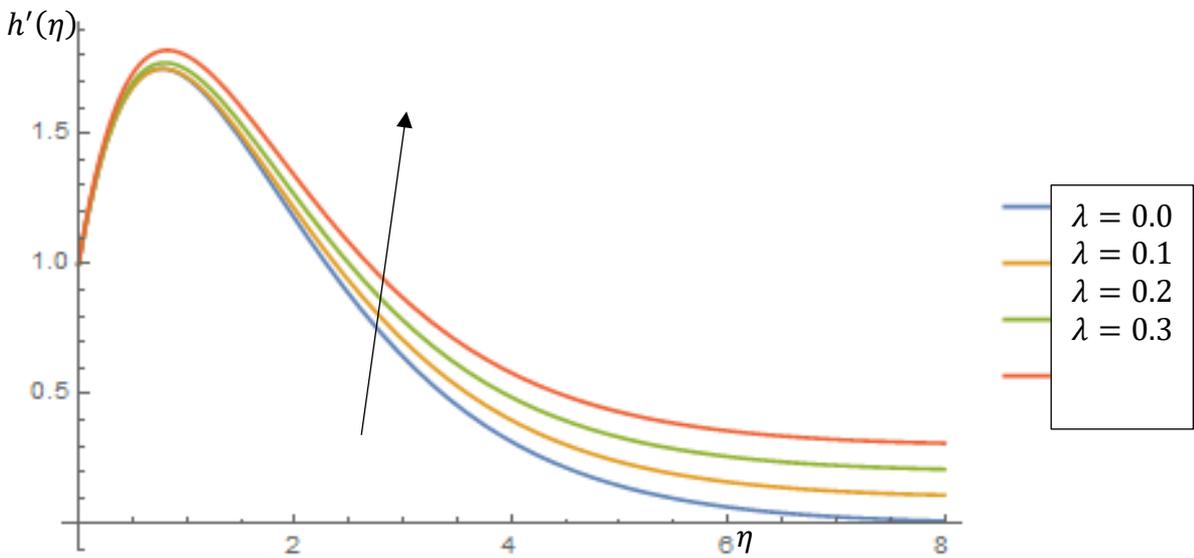


Figure 11: Impact of the stretching sheet parameter λ , on velocity with $M = R = 0.1, Gr = 5.0, Pr = 1.1$

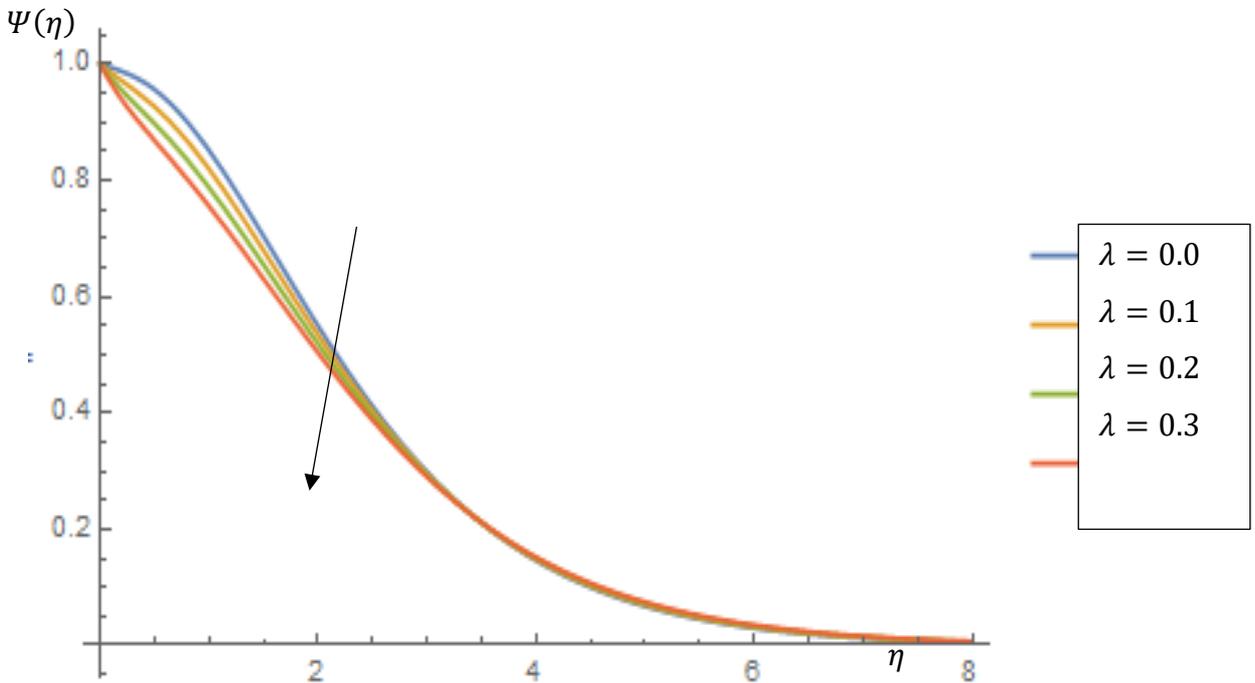


Figure 12: Impact of the stretching sheet parameter λ , on Concentration with $R = Nt = Nb = \gamma = 0.1, Gr = 5.0, Pr = 1.1$.

5. Discussion of Results

The analysis of the fluid parameters involved in this study are presented as follows:

The vertical axis of the profiles with their corresponding legends are represented by $h'(\eta)$, $\theta(\eta)$, $\Psi(\eta)$ i.e., velocity, temperature and concentration while on the horizontal axis, we have η in each case. Meanwhile, the legends displayed the effect of each parameter on the system. Thus, such impacts on the flow are discussed as follows:

The effect of magnetic strength M , on nanofluid motion is illustrated in figure 1. The magnetic field produces an opposing strength termed Lorentz force that slows down boundary layer flow making the fluid movement to decrease as depicted in figure 1. In order to bring the momentum development under control, the magnetic field is needed. Thus, depth of the momentum hydrodynamic reduces as magnetic parameter intensifies.

Figures 2 and 3 captures how radiation R is affecting velocity and temperature distributions. Growing R , leads to enhancement in the rate of fluid passage. The reason being that high level of thermal energy which exists in form of heat generated through thermal radiation intensifies, the force (bond) binding the various components of the fluid particles together is weakened and in turn leading to the

increase in fluid motion as depicted in figure 2. Similarly, the introduction of radiation in the path of the fluid motion means enhancing thermal radiation in and around the boundary and this entails surge in the boundary sheet thickness resulting in fluid temperature increase. This is illustrated in figure 3.

Figure 4 depicts the influence of thermal Grashof number Gr , on velocity distribution. Physically, Gr , is expressed as proportion of buoyancy to viscous forces and increasing its value within the momentum boundary layer, leads to increment in the buoyancy forces resulting to increase in velocity close to the sheet. This is due to the fact that closer to the sheet which radiates thermal energy in form of heat, the fluid molecules acquire more kinetic energy which makes them to be involved in more activity.

Significance of variations of Prandtl parameter Pr , on temperature profile is well illustrated in figures 5. From this figure, it can be seen that higher values of Prandtl number Pr , reduces the temperature thereby reducing the thermal boundary layer thickness. To support this phenomenon, we state that heat is diffused on heated sheet all through to the stretching sheet surface

The evolution of Brownian motion Nb , on nano-particle concentration is given by figures 6. When the value of Nb intensifies, a rapid movement of nanoparticles commence and this leads to a rise in acceleration which adds more energy to the particles, thus creating a fall in diffusion gradient of nanoparticles. This can be explained as the enhancement in Nb imparts on motion of particles of the fluid, in a way that such particles quickly migrate from areas of greater concentration towards areas of lesser concentration because of the fact that the nano-particle concentration decreases due to random motion of the particles.

Figures 7 and 8 entails the result of thermophoresis strength Nt , on temperature and concentration distribution. The thermophoresis movement of nano-particles leads to the movement energy in flow region towards top of the stretching sheet. Thus, a decline in temperature gradient evolves as exhibited in figure 7. However, an increase in thermophoresis parameter Nt , leads to a rise in concentration of particles. This is because intensifying thermophoresis parameter causes quick movement with the particles of the fluid flow which in turn brings about excess thermal energy in form of heat energy. Thus, this accounts for the high uprising of concentration profile.

Evolution on impact of Sc , on concentration profile is displayed in figures 9. As a dimensionless number, Sc is expressed as a relation of momentum diffusivity (fluid viscosity) to mass diffusivity. However, when there is more fluid viscosity than heat energy and mass diffusivity, an intensification in the Sc occurs around boundary layer. Hence, a rise in Schmidt number lowers the concentration. The impact of rate of chemical reaction is illustrated by profile 10. Increasing rate of chemical reaction is associated with a decay process in the thickness of the concentration boundary layer thereby causing decrease in its distribution.

The influence of stretching surface λ , on $h'(\eta)$ and $\Psi(\eta)$ distributions are captured in figures 11 and 12. Increasing the stretching sheet parameter λ is

accompanied with a rise in velocity shown in figures 11. However, this may be attributed to the fact that the stretching sheet velocity is less than the free stream velocity. Thus, the fluid movement tends towards the free stream region near the surface of the plate. Also, from figure 12, the trend changes because enhancing λ , lowers, $\Psi(\eta)$. Hence, this accounts for decrease in concentration, because of the fact that particles are in constant migration.

6. Conclusion

In this work presently, we have made some important findings concerning the motion of fluid, temperature as well as concentration of nanoparticles in nanofluid. However, the system of equations describing the flow are cracked through analytical means through the application of asymptotic series technique with MATHEMATICA.

We also considered the impact of flow constraints and conclude as follows:

1. The impact of magnetic strength parameter on the fluid flow indicates a decline in velocity as the parameter increases while the reverse is the case when thermal Grashof, radiation and stretching factors rise.
2. As the values of Prandtl, Schmidt and stretching numbers upsurge, the temperature and concentration distributions decrease.
3. The effect of radiation and thermophoresis factors is observed as the intensification leads to a reduction in both the temperature and concentration of nanoparticles.
4. The concentration distributions decreases as the parameters for Brownian motion and rate of chemical reaction increase.

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Nomenclature

(u, v)	the velocity components
μ	viscosity coefficient.
ρ	density of fluid.
σ	electrical conductivity of fluid.
M	Magnetic field parameter.
T	fluid temperature.
k	fluid heat conductivity term.
C_p	heat capacity at constant pressure.
q_r	heat flux radiative.
C	Concentration.
τ	the proportion of heat capacity of nanofluid to that of base fluid.
$(\rho C)_p$	nanofluid heat capacities.
$(\rho C)_f$	base fluid heat capacities.

D_B	Brownian motion coefficient.
D_T	thermophoretic diffusion coefficient.
T_∞	ambient fluid temperature.
C_∞	ambient fluid concentration.
Pr	Prandtl factor.
Sc	Schmidt parameter.
Nt	thermophoresis number.
Nb	Brownian motion factor.
λ	Stretching sheet velocity parameter.
R	thermal radiation factor.
Gr	convective thermal Grashof factor.
Gm	convective solutal Grashof factor.

Statement Interest

To the best of the authors knowledge, there's an absence of conflict of interest.

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